# EXPLORING POSSIBLE HIGH FUSION POWER REGIMES WITH THE IFS-PPPL MODEL

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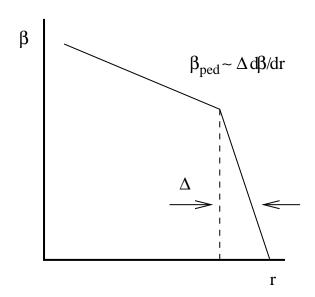
Some of this discussed by D. Meade (PPPL) at Workshop on Burning Plasma Sciences, General Atomics, May, 2001

Original: Workshop on Burning Plasma Sciences, December, 2000

Extension of paper presented at 1999 Snowmass workshop, http://www.ap.columbia.edu/SMproceedings see also PPPL-3360 (1999): http://www.pppl.gov/pub\_report/1999/PPPL-3360-abs.html

# Edge pedestal scalings very uncertain, but most favor higher-field designs with stronger shaping...

• Wide range of theory & expt. evidence:  $\Delta/R \propto \rho_{*\theta}$  (JT-60U, JET),  $\rho_{*\theta}^{2/3-1/2}$ ,  $\beta_{pol}^{1/2}\rho_{*}^{0}$  (very interesting DIII-D evidence of a second stable edge, which would have a more favorable scaling to reactors)



- ullet Making two assumptions (and use Uckan formula for  $q_{95}RI_p/(Ba^2)$ ):
  - 1. Width  $\Delta \propto \sqrt{\epsilon} \rho_{\theta} \propto \rho q/(\kappa \sqrt{\epsilon})$  (scaling preferred by two largest tokamaks)
  - 2. stability limit  $\partial \beta/\partial r \propto [1+\kappa^2(1+10\delta^2)]/Rq^2$  (rough fit to JT-60U, Koide et.al., Phys. Plasmas 4, 1623 (1997), other expts.), get:

$$T_{ped} = C_0 \left(\frac{n_{Gr}}{n_{ped}}\right)^2 \left[\frac{1 + \kappa^2 (1 + 10\delta^2)}{[1 + \kappa^2 (1 + 2\delta^2 - 1.2\delta^3)]} \frac{(1 - (a/R)^2)^2}{(1.17 - 0.65a/R)}\right]^2 \frac{A_i R}{\kappa^2 a}$$

## JET data supports $\Delta \propto \rho_{banana}$ & $\partial \beta/\partial r \propto Rq^2$ model.

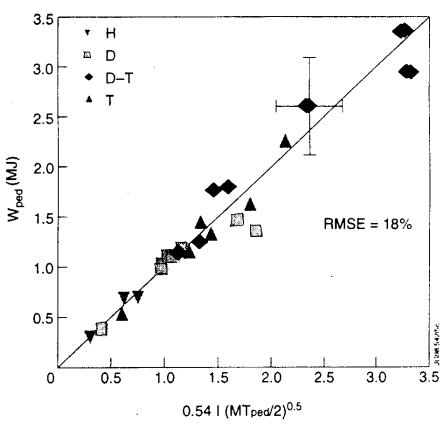


Fig. 4. Scaling of the stored energy in the pedestal (MJ) versus the fit 0.54 I  $(MT_{ped}/2)^{0.5}$ . The symbols are H=Hydrogen, D=Deuterium, D-T=50:50 D-T mixture and T=Tritium.

Cordey + JET Team, JAEA 198

JET data supports  $\triangle \propto \rho$  banana

+  $\frac{\partial \beta}{\partial r} \propto R_{g}^{2}$  model

# JT-60U showed the first evidence for the $\Delta \propto \rho_{banana}$ , $d\beta/dr \propto 1/(Rq^2)$ model. Also find a strong triangularity dependence.

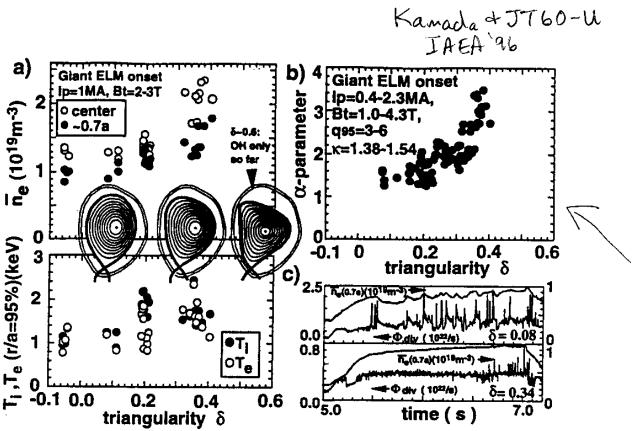


Fig. 1. a) and b): Increasing  $\overline{n}_e$  (center chord),  $\overline{n}_e$ (0.7a),  $T_e$ (r/a=95%),  $T_i$ (r/a=95%) and edge  $\alpha$ -parameter with increasing triangularity at onset of giant ELMs. c): Time traces of  $D_{\alpha}^{div}$  and  $\overline{n}_e$ (0.7a) for giant ELMs( $\delta$ =0.08) and grassy ELMs ( $\delta$ =0.34,  $\beta_p$ =2.4) with  $P_{NB}$ =20MW and  $I_p$ =0.6MA.

# Some of the new reactor designs may have significantly improved pedestal temperatures

Using this  $T_{ped}$  formula (with a  $\Delta \propto \rho_{\theta}$  assumption), and other pedestal scalings also, to scale from JET to some proposed reactor designs:

	R	а	В	$I_p$	$n_{ped}$	$\frac{n_{ped}}{n_{Gr}}$	$\frac{n_{ped}}{\langle n \rangle}$	$\kappa_{95}$	$\delta_{95}$	$T_{ped}$	$T_{ped}$	$T_{ped}$
	m	m	Т	MA	$10^{20}/m^3$	r•G1	(10)			keV	keV	keV
										if $\Delta \propto  ho_{ heta} \sqrt{\epsilon}$	if $5\delta^2$	if $\Delta \propto \sqrt{Rq ho}$
JET-norm	2.92	0.91	2.35	2.55	0.4	0.40	$\sim 1$	1.61	.17	2.1	2.1	2.1
ITER-96	8.14	2.80	5.68	21.0	1.3	1.52	1	1.60	.24	0.20*	0.18*	1.5*
lower $n_{ped}$	8.14	2.80	5.68	21.0	0.6	0.70	.70	1.60	.24	$0.94^{*}$	$0.83^{*}$	4.2*
ITER-FEAT	6.20	2.00	5.30	15.1	0.58	0.48	.65	1.70	.33	2.9	2.1	7.4
FIRE	2.0	0.53	10.0	6.44	3.6	0.48	.65	1.77	.40	4.8	3.0	6.7

<sup>\*</sup> should add  $(nT)_{sol}/n_{ped}$  which could be as high as  $\sim 0.5$  keV.

Encouraging that even with the pessimistic pedestal scaling ( $\Delta \propto \rho_{\theta}$ ), it may be possible to get high pedestal temperatures by going to stronger plasma shaping, higher field, smaller size, and modest density peaking.

## **Sensitivity of Fusion Power to Some Assumptions**

#### **Baseline assumptions:**

IFS-PPL model for  $\chi_{i,e}$  modified with  $\Delta(R/L_{Tcrit})=2$  to roughly fit Dimits shift seen in gyrokinetic simulations.

$$\langle n_e \rangle/n_{\rm Greenwald}=0.74$$
. Modest density peaking,  $n_0/\langle n_e \rangle=1.18$ ,  $n_{ped}/\langle n_e \rangle=0.65$ .  $n(r)=(n_0-n_{ped})(1-(r/a)^2)^{0.5}+n_{ped}$ .

 $P_{aux}$  adjusted to keep  $P_{net} \geq 1.2 P_{99L \rightarrow H}$  = 30 MW for baseline FIRE, =57 MW for baseline ITER-FEAT.

	$n_0$	$n_{ped}$	$T_{ped}$	$P_{fusion}$	Q	$T_{i0}$	$oxed{P_{aux}}{\sf MW}$
FIRE baseline case	$\frac{10^{20}/m^3}{6.75}$	$\frac{10^{20}/m^3}{3.6}$	keV 4.8	264	620.0	keV 18.6	IVIVV
			_		020.0		11
$\downarrow T_{ped}$ 30%	6.75	3.6	3.4	142	9.7	15.3	14
flatten $n(r)$	3.60	3.6	4.8	117	22.0	21.7	5
original IFS-PPPL	6.75	3.6	4.8	155	13.0	12.9	11
original IFS-PPPL $\downarrow T_{ped}$ 30%	6.75	3.6	3.4	69	2.6	10.2	26
ITER-FEAT baseline case	1.09	0.58	2.9	192	5.8	18.3	32
$\downarrow T_{ped}$ 30%	1.09	0.58	2.0	111	2.4	15.5	45
ITER-FEAT with FIRE $T_{ped}$	1.09	0.58	4.8	381	816.0	23.5	0
<b>ITER-FEAT</b> with FIRE $T_{ped} \downarrow 30\%$	1.09	0.58	3.4	241	10.1	19.8	23

### **CAVEATS, IMPLICATIONS**

- Dimits shift  $\Delta(R/L_{Trit}) \neq constant$ , should depend on parameters. Core neoclassical  $E \times B$  shear ignored (gets weaker at smaller  $\rho_*$ ).
- Edge pedestal scalings very uncertain.
- $T_{pedestal} \propto (n_{Greenwald}/n_{ped})^2$  model has no explicit power dependence, is only a guideline limit for certain regimes (first-stability-limited type-I ELMs). Assumes  $P > P_{LH}$  threshold. Ignores power needed to sustain pedestal against neoclassical transport, residual edge turbulence, ELMs, etc. Exploring extensions to include  $\nu_*$  dependence of bootstrap current, ...
- ullet To study edge turbulence & transport barriers scalings, need flexibility to scan pedestal density over a wide range: high  $n_{Gr}$ , pellet injection, divertor pumping.
- Compact size and strong shaping of FIRE gives high  $n_{Gr}$  & improved edge stability & high  $T_{pedestal}$  potential. Lower bound on  $n_{ped}$  needed for divertor survival appears to be easily satisified in FIRE.

### MORE CAVEATS, FUTURE WORK

Many caveats, contradictory theories, contradictory experiments:

- edge very complicated, range of theories, most have width  $\Delta \propto \rho^{2/3-1}$  .
- largest machines (JT-60U, JET) support "standard" model of width  $\Delta \propto \rho$  and gradient near the ideal MHD limit
- others (DIII-D) support  $\triangle$  independent of  $\rho$  and/or in second stability (boostrap current in pedestal region important in DIII-D?). C-MOD EDA differs from ELMy behaviour on other machines, Neutrals important in C-MOD?
- Useful cross-machine database being developed (Sugihara et.al., EPS99, ITER H-mode Edge Pedestal Expert Group Meeting, March 2000). (Sugihara uses different scaling  $dp/dr \propto (1+9.26\delta^{3.4})$ .)
- Detailed edge turbulence simulations rapidly becoming more realistic (Xu and Cohen (LLNL), Rogers and Drake (U. Md.), Scott, Jenko, Zeiler et.al. (Garching))
- Even with pessimistic  $\Delta \propto \rho$  model, newer reactor designs get significantly improved pedestal temperatures by  $\uparrow$  field, triangularity, and elongation (which increase Greenwald density and edge stability), and by assuming a modest density peaking

## May 2001 Addendum

- H-mode expts give evidence of multiple regimes: ELM-free, ELMY, Type-I, -II, -III, EDA. Different experiments show different scalings for pedestal width and height.
- Different physics may be setting limits in various regimes: The model presented here (pedestal width  $\Delta \propto \rho$  model with a first-stability beta limit) may be applicable in only certain regimes.
- In other regimes the edge bootstrap current may lower magnetic shear enough to lower the first stability boundary (Sugihara, EPS 1999) or even to access 2cd stability (as DIII-D expts and analysis by Osborne, Miller, et.al. suggest). However, if the edge bootstratp current gets too strong it may trigger a peeling mode (as Wilson, Snyder, etc. are studying). Studying improved mixed-regime models with Onjun, Bateman, Kritz (Lehigh).
- Hopefully these uncertainties can be reduced with the new edge database and comprehensive edge turbulence/stability simulations.